Exact solution for the diffusion controlled rate into a pair of reacting sinks^{a)}

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Recently there has been interest in determining the concentration dependence of diffusion controlled reactions. $^{1-3}$ The concentration dependence physically arises from the competition for the diffusing species among the reactive sites. This competition has the effect of reducing the effective rate coefficient of an assembly of reactive sites below the value expected from N independent sites, $k_0 = N(4\pi DR)$. Here D is the diffusion coefficient of the reacting species and R is the radius of the reaction center.

In this note, we present an exact solution for the case of two fixed centers N=2 in order to provide a basis for comparison to the more general theories¹ which, however, introduce approximations to deal with the case of greater interest, of large N.

The problem we solve is

$$D\nabla^2 n(\mathbf{r}) = 0, \tag{1}$$

where $n(\mathbf{r})$ is the concentration of the reacting species at point \mathbf{r} , subject to the boundary conditions that $n(\mathbf{r})$ - n_0 at large distances from the reaction center pair and that $n(\mathbf{r})$ vanishes at the boundary of the two centers that

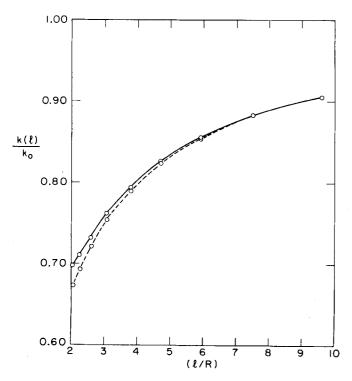


FIG. 1. Reduced diffusion controlled rate coefficient vs separation of reacting sinks. Solid line, exact result, Eq. (3); broken line, approximate result, Eq. (4).

are separated by a distance l, assigned to lie on the z axis. The solution, which is identical to Laplace's equation for the electrostatic potential surrounding two conducting spheres, is found in terms of bispherical coordinates⁴:

$$n(\mathbf{r})/n_0 = 1 - \left[2(\cosh \mu - \cos \eta) \right]^{1/2}$$

$$\times \sum_{m=0}^{\infty} \frac{\exp\left[- \left(m + \frac{1}{2} \right) \mu_0 \right]}{\cosh\left[\left(m + \frac{1}{2} \right) \mu_0 \right]}$$
(2)

$$\times \cosh[(m+\frac{1}{2})\mu]P_m(\cos\eta),$$

where the coordinates (μ, η, ϕ) are related to the cartesian coordinates (x, y, z) as given in Ref. 4. The rate coefficient is determined by computing the specific flux "disappearing" into both centers. The result is

$$\frac{k(l)}{k_0} = 1 - \sum_{m=2}^{\infty} (-1)^m \frac{\sinh \mu_0}{\sinh m \mu_0},\tag{3}$$

where the parameter μ_0 is determined by $\cosh \mu_0 = l/2R$. In the limit of infinite separation $l/2R \to \infty$ ($\mu_0 \to \infty$), the result is $k \to k_0$. At contact $l \to 2R(\mu_0 \to 0)$ the result is $k(2R) = (\ln 2)k_0 = 0.69 \ k_0$. The result $k(l)/k_0$ versus separation (l/R) is plotted in Fig. 1.

Included in Fig. 1 is the approximate result⁵

$$[k(l)/k_0]_{approx} = [1 + (R/l)]^{-1}$$
(4)

which agrees with Eq. (3) in the large separation limit where $l + R \exp(\mu_0)$. At contact the approximate expression gives the result $[k(2R)/k_0]_{approx} = \frac{2}{3}$, which is remarkably close to the exact result when one considers that the approximate method is essentially based on the assumption of large separation. This fortunate agreement suggests that the large separation approximation employed in theories of the concentration dependence of the diffusion controlled rate coefficient may not be too severe.

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⁵J. M. Deutch, B. U. Felderhof, and M. J. Saxton, J. Chem. Phys. **64**, 4559 (1976), and references cited therein.

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²D. Peak, K. Pearlman, and P. J. Wantuck, J. Chem. Phys. **65**, 5538 (1976).

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⁴P. M. Morse and H. Feshbach, Methods of Theoretical Physics, Part II (McGraw-Hill, New York, 1953), page 1298 ff.