

Lecture Notes on Fluid Dynamics

(1.63J/2.21J)

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2007 Spring

2-4-apile.tex

References:

Landau & Lifshitz, *Fluid Mechanics*, 1959.

H. Huppert : (JFM) (For modelling lava flow.)

2.4.a. Release of a fluid pile from a horizontal plane

On a horizontal plane $\theta = 0$. Eq (2.2.22) reduces to

$$\frac{\partial h}{\partial t} = \alpha \frac{\partial}{\partial x} \left(h^3 \frac{\partial h}{\partial x} \right) \quad (2.4.a.1)$$

where $\alpha = \rho g / 3\mu$. This is a nonlinear diffusion equation where the diffusivity ah^3 increases with the unknown h .

Consider the release of a fluid pile with the initial volume Q confined in a small region. Note that

$$\frac{\partial}{\partial t} \int_{-\infty}^{\infty} h(x, t) dx = \alpha \left[h^3 \frac{\partial h}{\partial x} \right]_{-\infty}^{\infty} = 0$$

Thus

$$\int_{-\infty}^{\infty} h(x, t) dx = \int_{-\infty}^{\infty} h(x, 0) dx = Q \quad (2.4.a.2)$$

i.e., the total mass is constant. The initial-value problem governed by (2.4.a.1) and (2.4.a.2) can be solved by the method of similarity, which reduces the PDE to an ODE (Landau & Lifshitz).

2.4.a.1. The similarity method

Consider the one-parameter transformation:

$$h = \lambda^a h', \quad x = \lambda^b x', \quad t = \lambda^c t' \quad (2.4.a.3)$$

where (a, b, c) will be chosen to leave the initial value problem (PDE and initial condition) unchanged (invariant). Substituting (2.4.a.3) into (2.4.a.1),

$$\lambda^{a-c} \frac{\partial h'}{\partial t'} = \lambda^{4a-2b} \alpha \frac{\partial}{\partial x'} \left(h'^3 \frac{\partial h'}{\partial x'} \right) \quad (2.4.a.4)$$

and then in (2.4.a.2),

$$\lambda^{a+b} \int_{-\infty}^{\infty} h'(x, 0) dx' = Q \quad (2.4.a.5)$$

we find the conditions for invariance,

$$a - c = 4a - 2b; \quad a + b = 0; \quad i.e., \quad b = -a, \quad c = -5a; \quad (2.4.a.6)$$

implying also

$$h = \lambda^a h', \quad x = \lambda^{-a} x', \quad t = \lambda^{-5a} t' \quad (2.4.a.7)$$

This suggests the following combination of variables:

$$\xi = \frac{x}{(At)^{1/5}}, \quad f(\xi) = h(Bt)^{1/5}, \quad (2.4.a.8)$$

ξ is called a *similarity* variable. In the space-time plane, at all points on the curve of constant $\xi = x/(Bt)^{1/5}$, $f = h(Bt)^{1/5}$ is the same. Thus for a fixed ξ , $h(Bt)^{1/5}$ is fixed, implying that $h(Bt)^{1/5}$ is a function of ξ , i.e.,

$$h = \frac{f(\xi)}{(Bt)^{1/5}} \quad (2.4.a.9)$$

Physically the maximum of h decays in time as $t^{-1/5}$, while the pile spreads in time as $t^{1/5}$.

Substituting (2.4.a.9) into (2.4.a.1), and using the facts,

$$\frac{\partial h}{\partial t} = -\frac{1}{5} \frac{1}{B^{1/5}} \frac{f}{t^{6/5}} - \frac{1}{5} \frac{\xi f'}{B^{1/5} t^{6/5}} \quad \frac{\partial h}{\partial x} = \frac{1}{B^{1/5} t^{1/5}} \frac{f'}{A^{1/5} t^{1/5}}$$

we find

$$-\frac{1}{5} \frac{1}{B^{1/5} t^{6/5}} \left(f + \xi \frac{df}{d\xi} \right) = \frac{\alpha}{B^{4/5} t^{4/5} A^{2/5} t^{2/5}} \frac{d}{d\xi} \left(f^3 \frac{df}{d\xi} \right)$$

From the initial condition, we get

$$\left(\frac{At}{Bt} \right)^{1/5} \int_{-\infty}^{\infty} f(\xi) d\xi = Q$$

For simplicity let us set

$$A^{2/5} B^{3/5} = \alpha \quad (2.4.a.10)$$

and

$$\left(\frac{A}{B} \right)^{1/5} = Q, \quad (2.4.a.11)$$

We then get an ODE:

$$5 \frac{d}{d\xi} \left(f^3 \frac{df}{d\xi} \right) + \frac{d}{d\xi} (\xi f) = 0 \quad (2.4.a.12)$$

subject to

$$\int_{-\infty}^{\infty} f(\xi) d\xi = 1. \quad (2.4.a.13)$$

at $t \rightarrow 0$. Solving (2.4.a.10) and (2.4.a.11) we get

$$A = \alpha Q^3, \quad B = \frac{\alpha}{Q^2} \quad (2.4.a.14)$$

2.4.a.2. Solution

Integrating (2.4.a.12),

$$5f^3 \frac{df}{d\xi} + \xi f = 0$$

The integration constant is zero because of symmetry at $x = 0$ so that $f'(0) = 0$. Now

$$5f^2 \frac{df}{d\xi} + \xi = 0, \quad \text{implying} \quad \frac{5}{3} df^3 + \xi d\xi = 0$$

Integrating,

$$f^3 = \frac{3}{5} \frac{\xi_0^2 - \xi^2}{2},$$

where $f(\xi_0) = 0$, we get

$$f = \left[\frac{3}{10} (\xi_0^2 - \xi^2) \right]^{1/3}, \quad -\xi_0 < \xi < \xi_0; \quad = 0, \quad \text{otherwise.} \quad (2.4.a.15)$$

To find ξ_0 we use (2.4.a.13),

$$\int_{-\xi_0}^{\xi_0} f(\xi) d\xi = 1.$$

or

$$\int_{-\xi}^{\xi_0} [\xi_0^2 - \xi^2]^{1/3} d\xi = \left(\frac{10}{3} \right)^{1/3}$$

This determines ξ_0 . The rest is algebra.

Let $\zeta = (\xi/\xi_0)^2$, then

$$d\zeta = 2\xi d\xi / \xi_0^2, \quad d\xi = \frac{\xi_0}{2} \frac{d\zeta}{\zeta^{1/2}}$$

$$\int_{-\xi}^{\xi_0} [\xi_0^2 - \xi^2]^{1/3} d\xi = (\xi_0)^{5/3} \int_0^1 (1 - \zeta)^{1/3} \zeta^{-1/2} d\zeta = (\xi_0)^{5/3} \frac{2\sqrt{\pi} \Gamma\left(\frac{1}{3}\right)}{5 \Gamma\left(\frac{5}{6}\right)}$$

where $\Gamma(z)$ is the Gamma function, hence

$$\xi_0 = \frac{x_0}{(\alpha Q^3 t)^{1/5}} = \left(\frac{10}{3} \right)^{1/5} \left(\frac{5}{2\sqrt{\pi}} \frac{\Gamma\left(\frac{5}{6}\right)}{\Gamma\left(\frac{1}{3}\right)} \right)^{3/5} \quad (2.4.a.16)$$

The maximum width is

$$x_0 = \left(\frac{10}{3} \right)^{1/5} \left(\frac{5}{2\sqrt{\pi}} \frac{\Gamma\left(\frac{5}{6}\right)}{\Gamma\left(\frac{1}{3}\right)} \right)^{3/5} (\alpha Q^3 t)^{1/5} \quad (2.4.a.17)$$

The maximum depth is

$$h(0, t) = \frac{f(0)}{(Bt)^{1/5}} = \left[\frac{3\xi_0^2}{10} \right]^{1/3} \left(\frac{Q^2}{\alpha t} \right)^{1/5} \quad (2.4.a.18)$$